24

## 9 Study of $\pi K$ atoms

C. Amsler, A. Benelli, and J. Rochet

(DIRAC Collaboration)

The main initial goal of the DIRAC experiment was to produce  $\pi^+\pi^-$  atoms (pionium) and to measure the  $\pi\pi$  scattering lengths  $a_0$  and  $a_2$  (isospins 0 and 2) [1]. The difference in the scattering lengths was measured,

$$|a_0 - a_2| = 0.2533 \pm 0.0109 \ m_{\pi}^{-1}$$

based on the observation of 21'227  $\pi^+\pi^-$  atoms. This value corresponds to a mean life of 3.15  $\pm$  0.28 fs. Details can be found in ref. [2].

We joined the DIRAC collaboration in 2007 with the goal to search for  $\pi K$  atoms and to measure the  $\pi K$  scattering lengths  $a_{1/2}$  and  $a_{3/2}$  (isospins 1/2 and 3/2).

An artist's view of the DIRAC spectrometer and a photograph of the experiment in the CERN-PS-South Hall is shown in figures 9.1 and 9.2, respectively. The 24 GeV/c proton beam passes through a thin target (such as 100  $\mu m$  Ni). The secondary particles traverse a scintillation fiber detector (SFD). The secondary hadrons produced in target are analyzed in a double-arm magnetic spectrometer measuring the momentum vectors of two oppositely charged hadrons. The spectrometer is slightly tilted up-

wards with respect to the proton beam. Electrons and positrons are vetoed by  $N_2$ -Čerenkov detectors (N) and muons by their signals in scintillation counters (Mu) behind steel absorbers. The preshower detector (PSh) provides an additional electron/hadron separation to the  $N_2$ -Čerenkov. Details can be found in previous annual reports and in ref. [3].

In 2007 we modified the DIRAC setup to study  $\pi K$  atoms while at the same time collecting more data for  $\pi\pi$  atoms. In both arms the N<sub>2</sub>-Čerenkov counters were cut to leave space for heavy gas (C<sub>4</sub> F<sub>10</sub>) Cerenkov detectors (HG). They identify pions with an efficiency of more than 99% while not responding to kaons nor to (anti)protons. On the left side we installed an aerogel Čerenkov detector (Ae) to identify kaons and reject protons. The detector consists of two modules with refractive index 1.015 and one module with refractive index 1.008 to cover the high momentum range of kaons between 5.5 and 8 GeV/c. Our group has developed and built the  $37\ell$  aerogel Čerenkov counter [4] and the gas system for the C<sub>4</sub>F<sub>10</sub> counters [5].

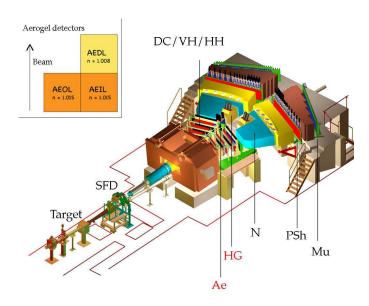


Fig. 9.1 - DIRAC spectrometer:

SFD: scintillating fibre detector, DC: drift chambers, VH (HH): vertical (horizontal) scintillating hodoscopes, PSh: preshower, Mu: muon counters, Ae: aerogel, and HG: heavy gas and N:  $N_2$ -Čerenkov counters. The inset shows the arrangement of the aerogel modules.



Fig. 9.2 – The DIRAC apparatus.

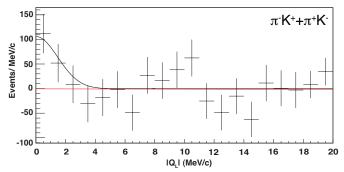


Fig. 9.3 – Distribution of the  $\pi K$  relative momentum in the c.m.s system after background subtraction. A Gaussian fit (solid line) illustrates the signal distribution.

Using only the downstream part of DIRAC for tracking we published the first evidence for  $\pi K$  atoms [6, 7]. A signal of  $173 \pm 54$  (3.2 $\sigma$ )  $\pi K$ -pairs was observed at very low relative momenta (typically  $|Q_L| < 3$  MeV/c in the c.m.s system), see Fig. 9.3. The evidence was strengthened by the simultaneous observation of interacting unbound Coulomb-pairs from which the fraction of expected bound pairs can be calculated. The latter is in excellent agreement with the number of observed atoms. This result corresponds to a lower limit on the mean life of  $\pi K$ -atoms of 0.8 fs in the 1s-state (confidence level of 90%), which translates to [6, 7]:

$$|a_{1/2} - a_{3/2}| < 0.58 \ m_{\pi}^{-1}.$$

More data was collected in 2008 - 2011 from which the total number of reconstructed  $\pi K$  atoms will be around 600. Assuming the central value as in 2007 the significance will increase to  $5.3\sigma$  corresponding to an uncertainty of about 25% on  $|a_{1/2} - a_{3/2}|$ .

In 2012 DIRAC also studied the energy splitting  $\Delta E$  between the pionium 2p- and 2s-states by mean of the

following method: pionium atoms produced in the target interact with an applied electric field and some of them leave the target in the 2*p*-state. The decay into  $\pi^0\pi^0$  is forbidden from p-states and hence the main decay process is the 2p - 1s radiative transition with subsequent annihilation from 1s into  $\pi^0\pi^0$ . Thus the mean life of the atom in the 2*p*-state is determined by radiative transition,  $\tau(2p) \simeq 12$  ps, which is much slower than annihilation from the 1s-states. We refer to these 2p-states as longlived  $2\pi$  atoms. One then measures  $\Delta E$  by observing the field dependence of the number of long-lived  $2\pi$  atoms breaking up into  $\pi^+\pi^-$  pairs in a second foil downstream of the production target. This measures the combination  $2a_0 + a_2$ . Combined with the measurements from the 1sstate, the  $\pi\pi$  scattering lengths  $a_0$  and  $a_2$  can be determined separately.

Data taking for the DIRAC experiment was completed at the end of 2012 and the apparatus was subsequently dismantled.

- B. Adeva et al. (DIRAC Collaboration),
  Phys. Lett. B 704 (2011) 24.
- [2] A. Benelli, EPJ Web of Conferences 37 (2012) 01011, doi: 10.1051/epjconf/20123701011.
- [3] B. Adeva *et al.* (DIRAC Collaboration), Nucl. Instrum. Methods in Phys. Res. **A 515** (2003) 467.
- [4] Y. Allkofer et al., Nucl. Instr. Meth. in Phys. Res. A 582 (2007) 497; Y. Allkofer et al., Nucl. Instr. Meth. in Phys. Res. A 595 (2008) 84; C. Amsler, Proc. of Science PoS EPS-HEP (2009) 078.
- [5] S. Horikawa et al.,Nucl. Instr. Meth. in Phys. Res. A 595 (2008) 212.
- [6] B. Adeva et al. (DIRAC Collaboration), Phys. Lett. B 674 (2009) 11.
- [7] Y. Allkofer, PhD Thesis, University of Zurich (2008).